

# Solar Wind and CMEs with the Space Weather Modeling Framework

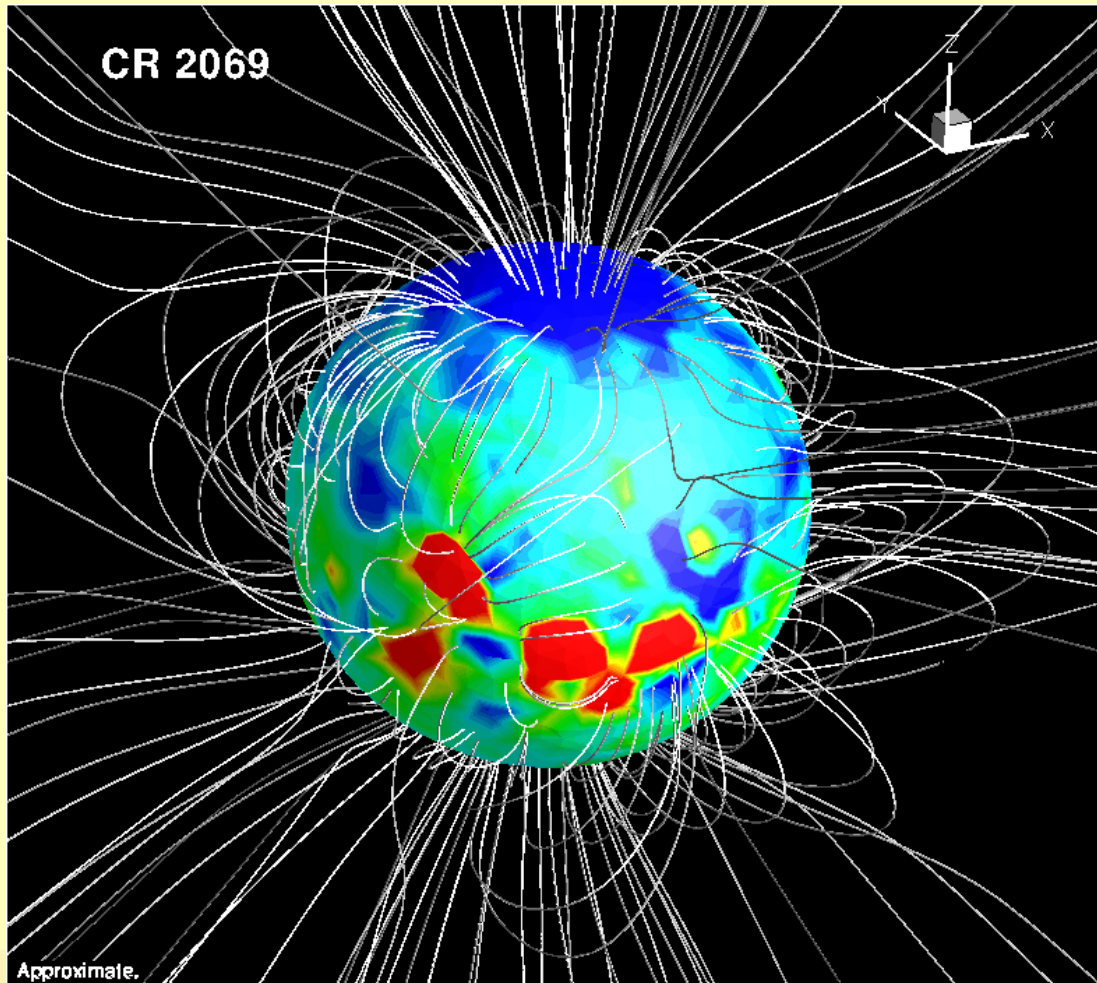
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University of Michigan



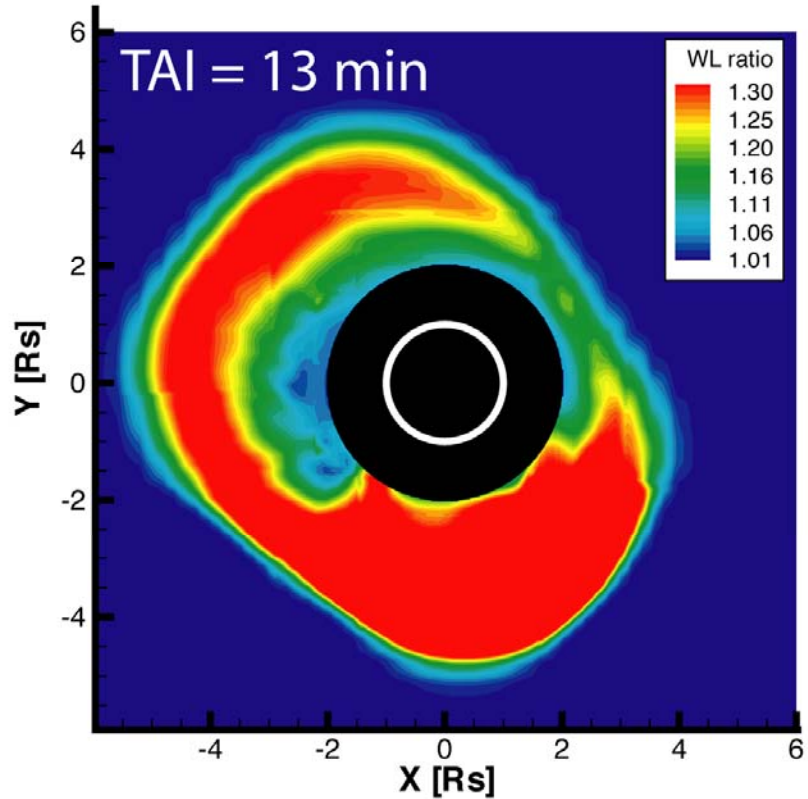
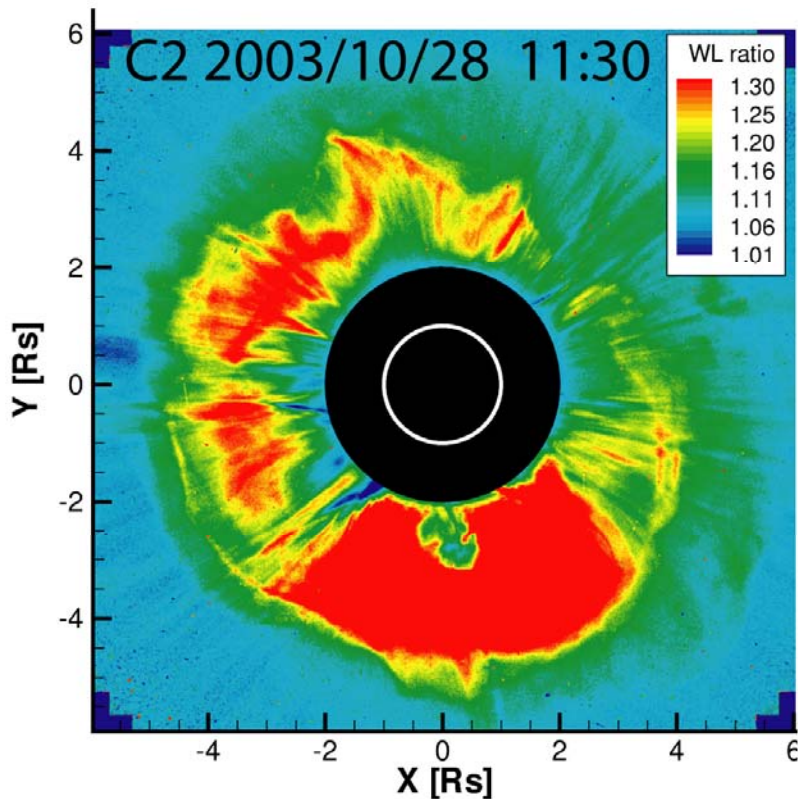
- M** Default Solar Corona model at CCMC
- M** New: Two-temperature corona model with Alfvén waves
- M** New: Lower Corona model with EUV images
- M** New: Flux Emergence in Convection Zone model
- M** Future plans

## Default Solar Corona Model at CCMC (O. Cohen et al.)



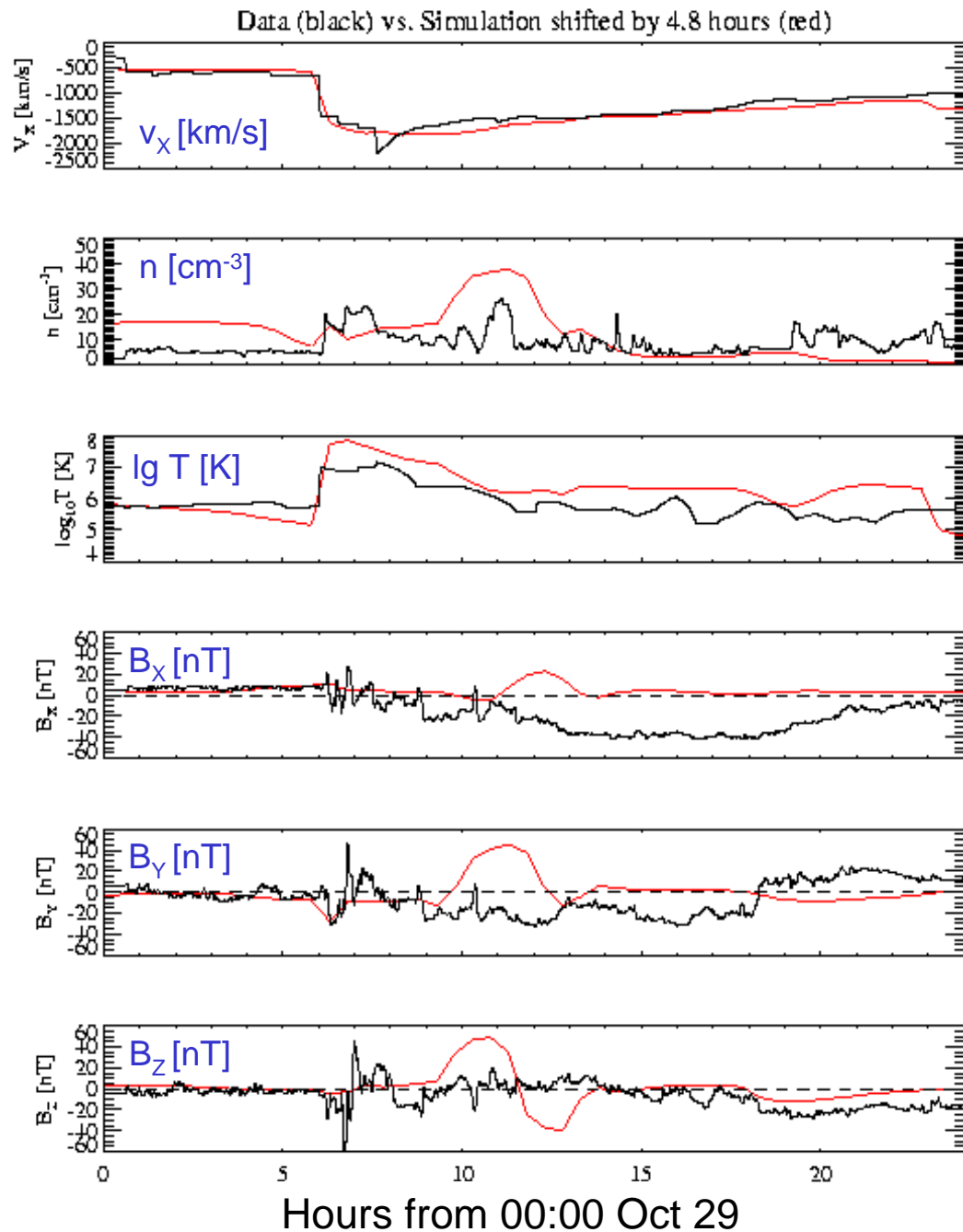
- M** Uses a spatially varying polytropic index  $\gamma$  to account for the coronal heating
- M** Uses the Wang-Sheeley-Arge model and Bernoulli integral along the field lines to determine  $\gamma$
- M** Designed to reproduce the observations at 1AU

# Quantitative Comparison with SOHO for the October 2003 Halloween Storm



- M** Out of equilibrium flux rope superposed on a steady state MHD corona
- M** Total Mass, Observed =  $1.5 \times 10^{16}$  gram Model =  $2.2 \times 10^{16}$  gram

# Observed vs. Simulated Solar Wind and IMF near Earth (10° off)



Simulated CME arrives 4.8 hours early.

Solar wind speed is almost perfect!

Density is within range.

Magnetic field variation has right amplitude but sign and phase differ.

## **M** Midterm review of CCHM (Comprehensive Corona and Heliosphere Model):

- 🌐 **Default Coronal model is based on a reduced, spatially varying adiabatic index for the heating and acceleration of the plasma.**
- 🌐 **In good agreement with the ambient solar wind observed at 1AU, but it is not self-consistent and distorts the physics, especially at CME shocks**

**M** In addition: Default CME is not self-consistent, but a superposed flux rope model.



**M** Use polytropic index  $\gamma = 5/3$

**M** Electrons are no longer heated by the ions beyond  $1.5R_s$ , where electron-ion collisions is infrequent

**M** Ion heating (by Alfvén waves) is redistributed by ion heat conduction (35 times smaller than electron heat conduction)

$$\frac{\partial p_i}{\partial t} + \nabla \cdot (p_i \mathbf{u}) + (\gamma - 1)p_i \nabla \cdot \mathbf{u} = (\gamma - 1) [Q_i - \nabla \cdot \mathbf{q}_i + \lambda_{ei}(T_e - T_i)],$$

$$\frac{\partial p_e}{\partial t} + \nabla \cdot (p_e \mathbf{u}) + (\gamma - 1)p_e \nabla \cdot \mathbf{u} = (\gamma - 1) [ -\nabla \cdot \mathbf{q}_e + \lambda_{ei}(T_i - T_e)],$$

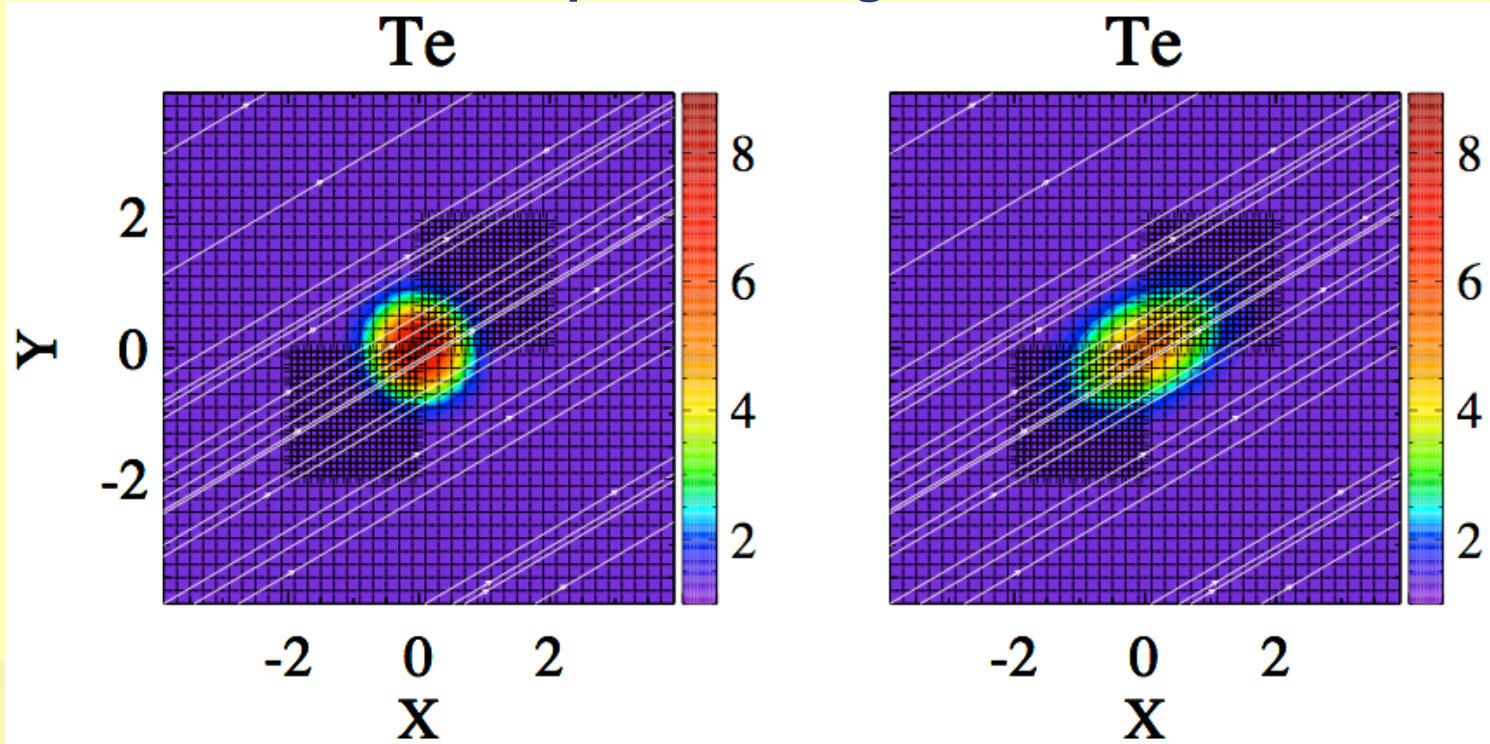
$$\mathbf{q}_i = -\kappa_i T_i^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_i, \quad \mathbf{q}_e = -\kappa_e T_e^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_e$$

**M** No collisionless heat conduction and additional electron heating yet

**M** Verification of the new collisional heat conduction implementation

$$\frac{\partial E_e}{\partial t} = \dots + \nabla \cdot \left( \kappa_e T_e^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_e \right)$$

**M** Diffusion of a Gaussian profile along the field lines

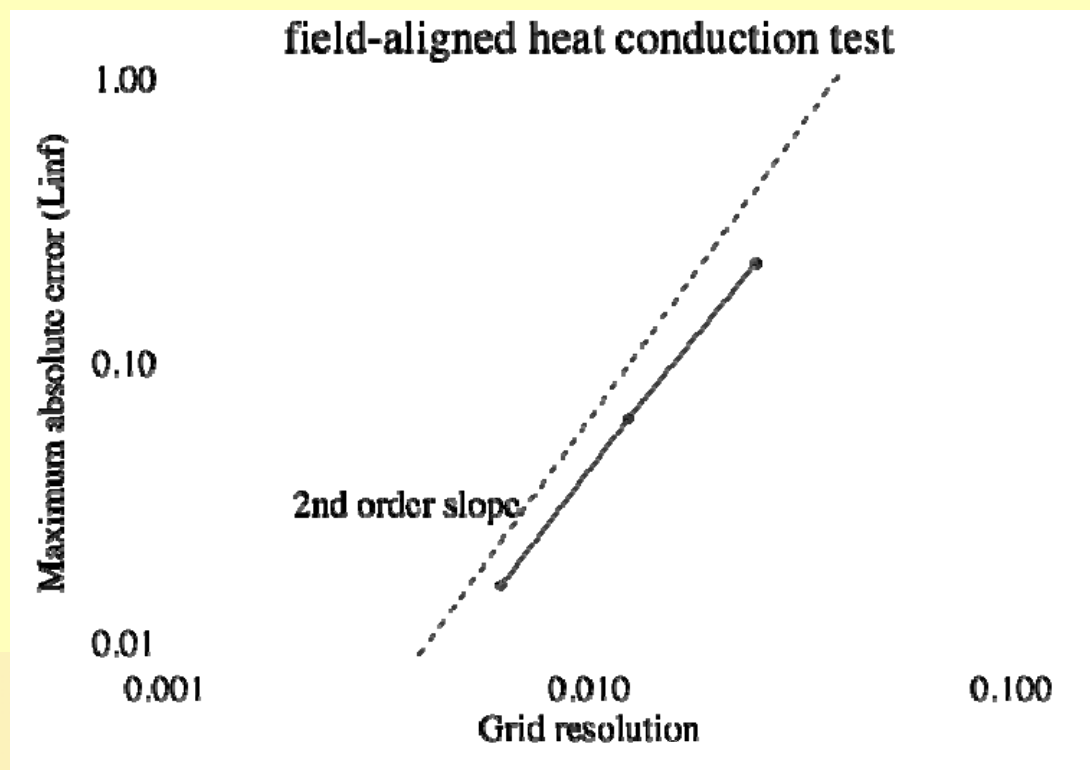




## M Verification of the new collisional heat conduction implementation

$$\frac{\partial E_e}{\partial t} = \dots + \nabla \cdot \left( \kappa_e T_e^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_e \right)$$

## M Diffusion of a Gaussian profile along the field lines



# WKB Alfvén Waves



- M** Hinode observations suggest that the Alfvén wave energy input is sufficient for heating and acceleration (e.g. de Pontieu et al. 2007).
- M** Wind acceleration: work done by wave pressure force
- M** Coronal heating: formulation of the Kolmogorov dissipation by Hollweg (1986)

$$\frac{\partial E_w^+}{\partial t} + \nabla \cdot [E_w^+ (\mathbf{u} + \mathbf{u}_A)] + p_w^+ \nabla \cdot \mathbf{u} = -Q^+,$$

$$\frac{\partial E_w^-}{\partial t} + \nabla \cdot [E_w^- (\mathbf{u} - \mathbf{u}_A)] + p_w^- \nabla \cdot \mathbf{u} = -Q^-,$$

$$\text{Ion heating } Q_i = Q^+ + Q^- = \frac{E_w^{+3/2}}{L\sqrt{\rho}} + \frac{E_w^{-3/2}}{L\sqrt{\rho}}, \quad L = C/\sqrt{B},$$

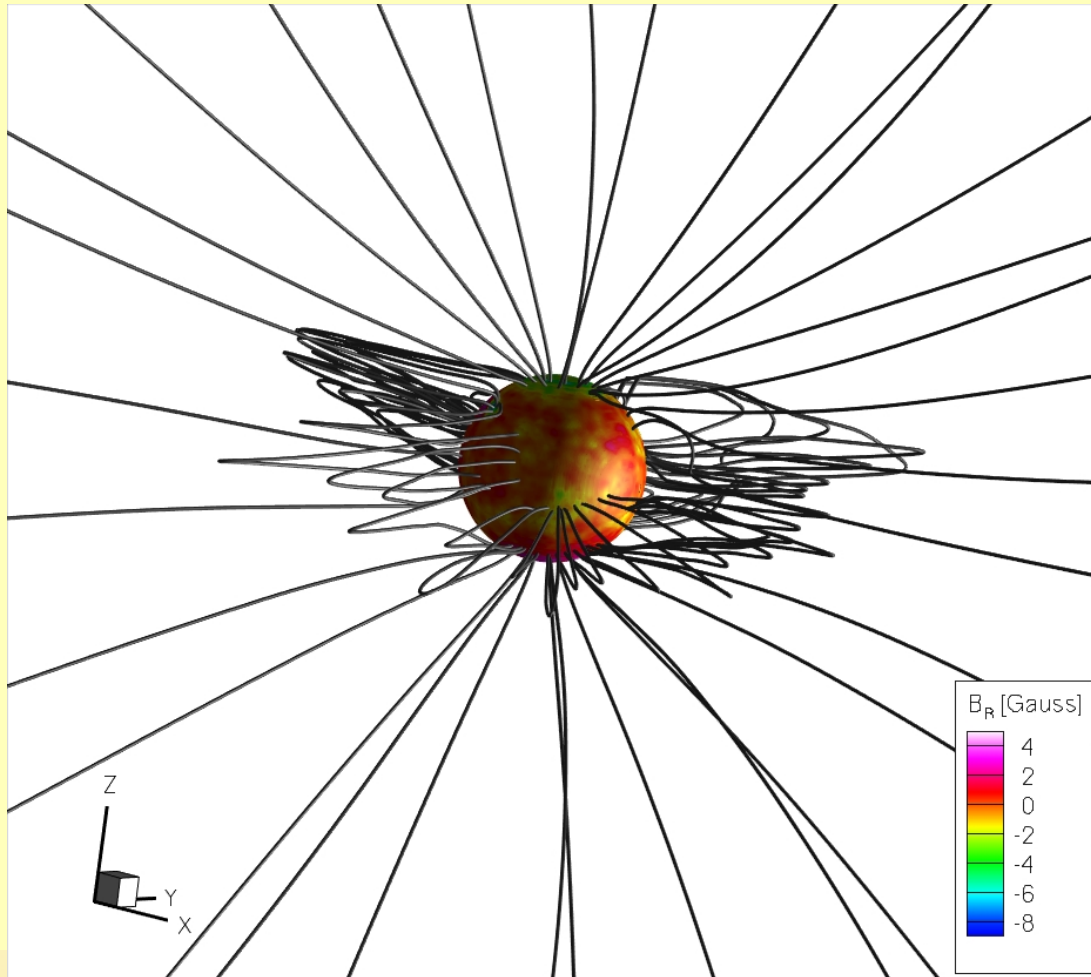
- M** Free parameter C in heating scale height L

- M** Magnetogram driven potential field extrapolation.
- M** Scale height of the heating and Alfvén wave dissipation.
- M** Specify Alfvén wave pressure at the solar base:
  - 🌍 Given solar wind expansion factor from PFSS model,
  - 🌍 Given velocity at 1AU based on WSA model,
  - 🌍 Impose energy conservation along the field lines to determine the Alfvén wave pressure at  $1R_S$ .

$$\frac{u_{\text{WSA}}^2}{2} = \frac{\gamma}{\gamma - 1} \frac{p_i + p_e}{\rho} - \frac{GM_{\odot}}{R_{\odot}} + R_{\odot}^2 \frac{|u_{A,r}| E_w}{(\rho u r^2)_{1\text{AU}} f_{\text{expansion}}}$$

- 🌍 In-going Alfvén waves at  $1R_S$  are absorbed.

## Result for CR2077



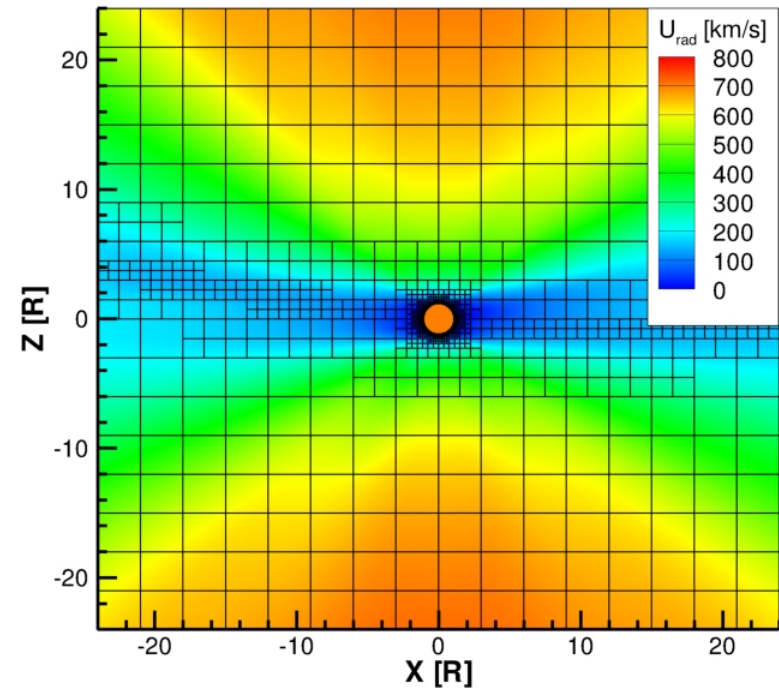
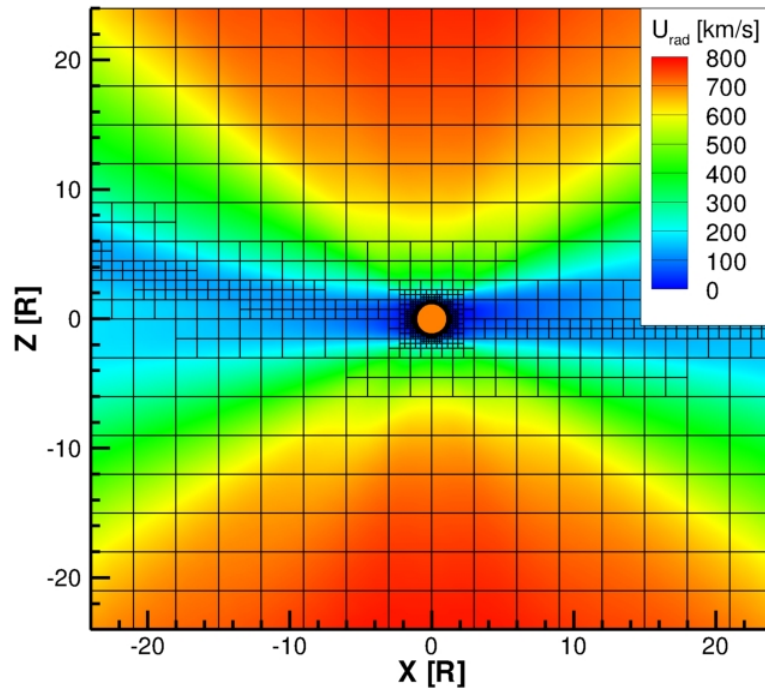
- M** SOHO-MDI magnetogram CR2077
- M** Calculated on a 3D cartesian grid with AMR
- M** Predicted rms velocities of the Alfvén perturbations near the sun is of the order of 30km/s

## 2T versus 1T model for CR2077



Radial velocity for 2 temperature model  
Maximum is 785 km/s

Radial velocity for 1 temperature model  
Maximum is 714 km/s

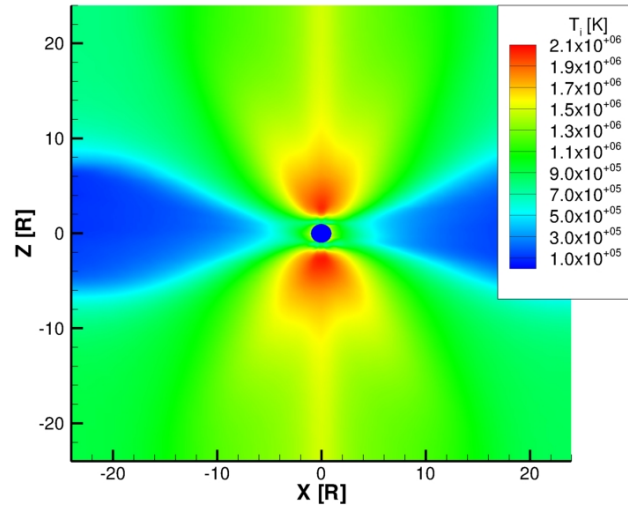


**M** The black squares represent the AMR blocks of 4x4x4 cells

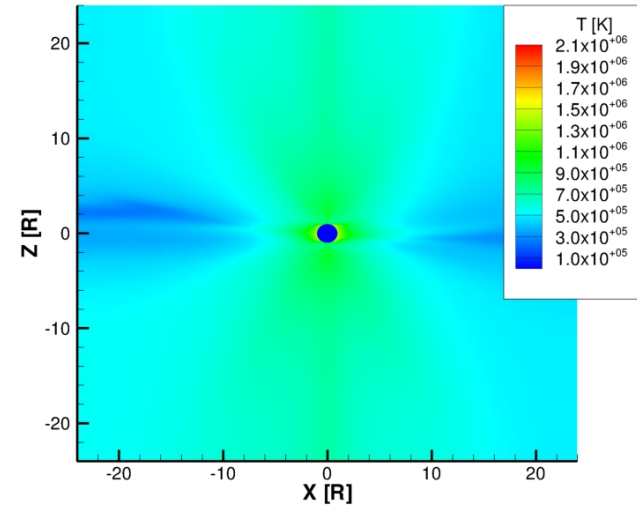
# 2T versus 1T model for CR2077



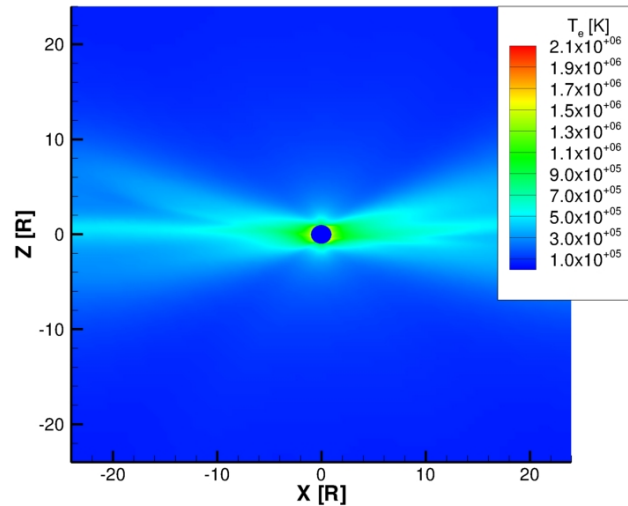
Ion temperature



Temperature



Electron temperature



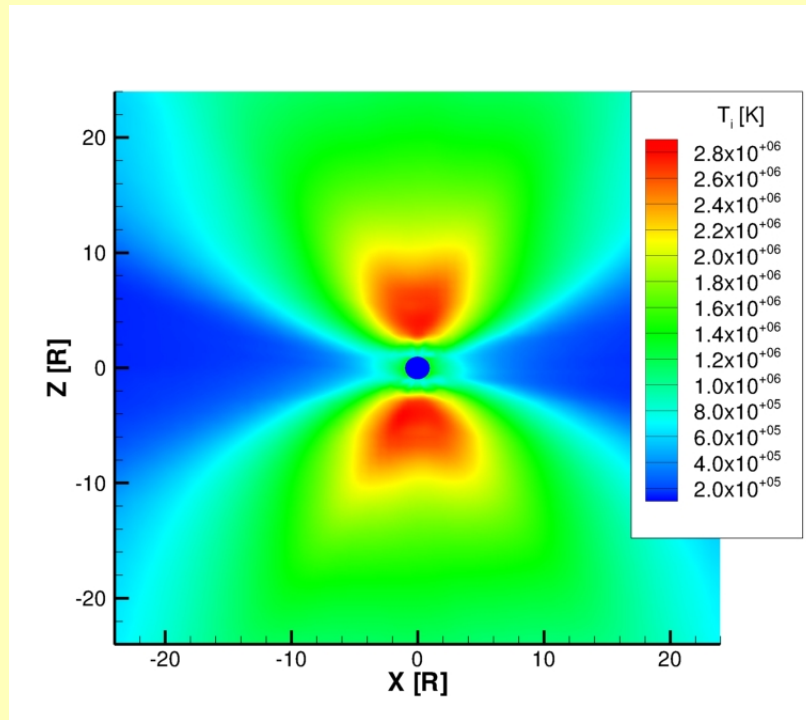
- M** In the 2T model, the ions are preferentially heated in the coronal hole.
- M** In the 1T model: The strong spatial redistribution of the ion heating by the electron heat conduction is unphysical.



# Importance ion heat conduction

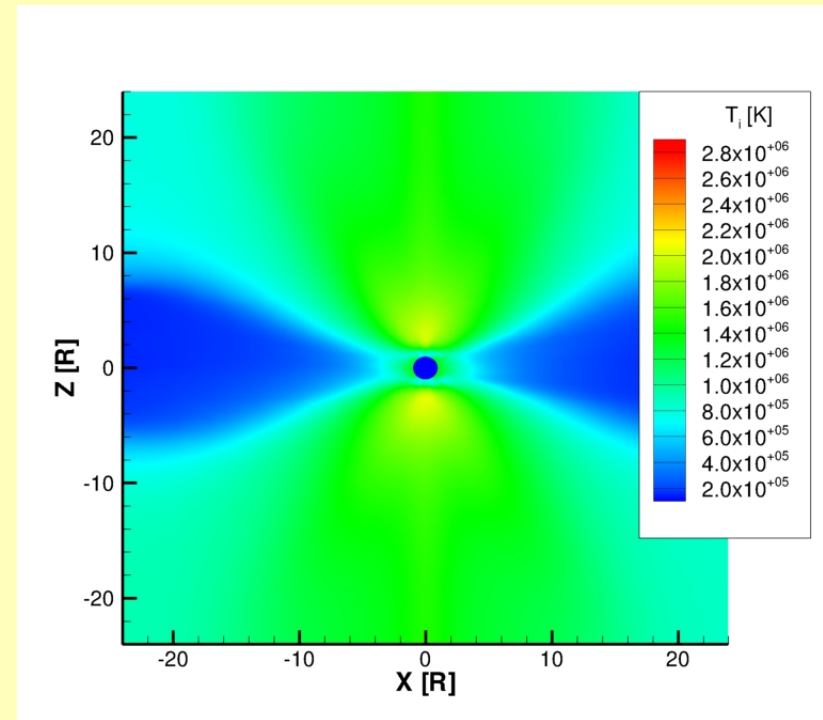


**M** Why using ion heat conduction if it is 35 times smaller than electron heat conduction ?



**M** 2T simulation without ion heat conduction

**M** Maximum ion temperature is 2.8MK



**M** 2T simulation with ion heat conduction

**M** Maximum ion temperature is 2.1MK

## **M Improved thermodynamics:**

- One temperature model up to  $1.5R_s$
- Electron heat conduction with artificial broadening of the transition region (similar to Lionello et al. 2009)
- Optically thin radiative losses
- Exponential heating function, in active regions transition to heating proportional to magnetic field strength

## **M Chromosphere boundary ( $T=20000\text{K}$ , $n_e=10^{12}\text{ cm}^{-3}$ )**

## **M Ability to synthesize line-of-sight images in the EUV and X-ray regime**

## **M Submitted to ApJ (C. Downs et al.)**

# EUV and X-ray LOS images for CR1913

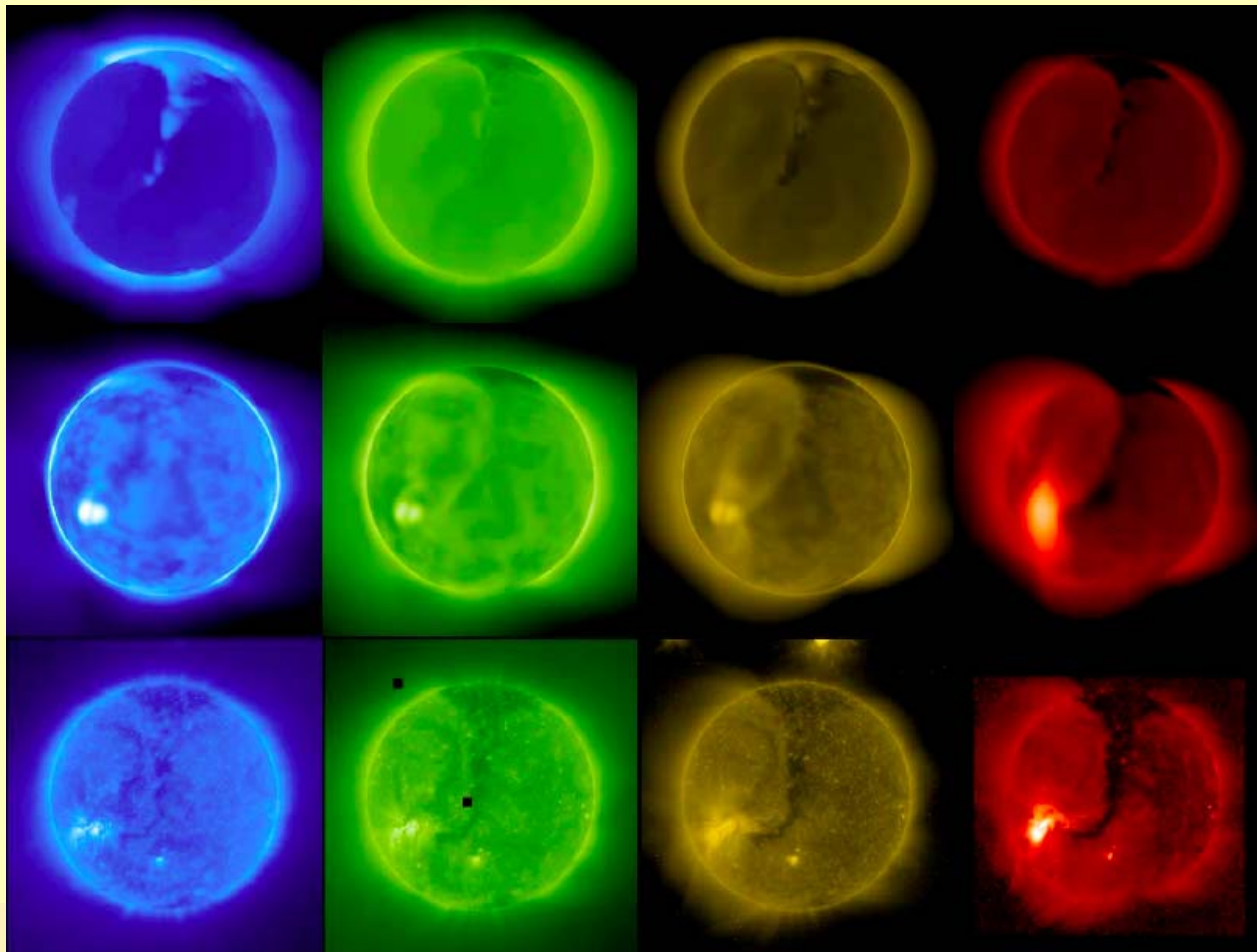


EIT 171Å

EIT 195Å

EIT 284Å

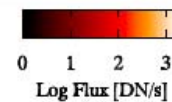
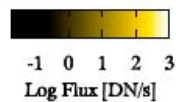
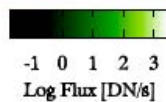
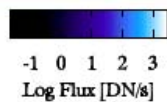
SXT AlMg



Old SC model synthesis

New LC model synthesis

Observation:  
Aug 27, 1997



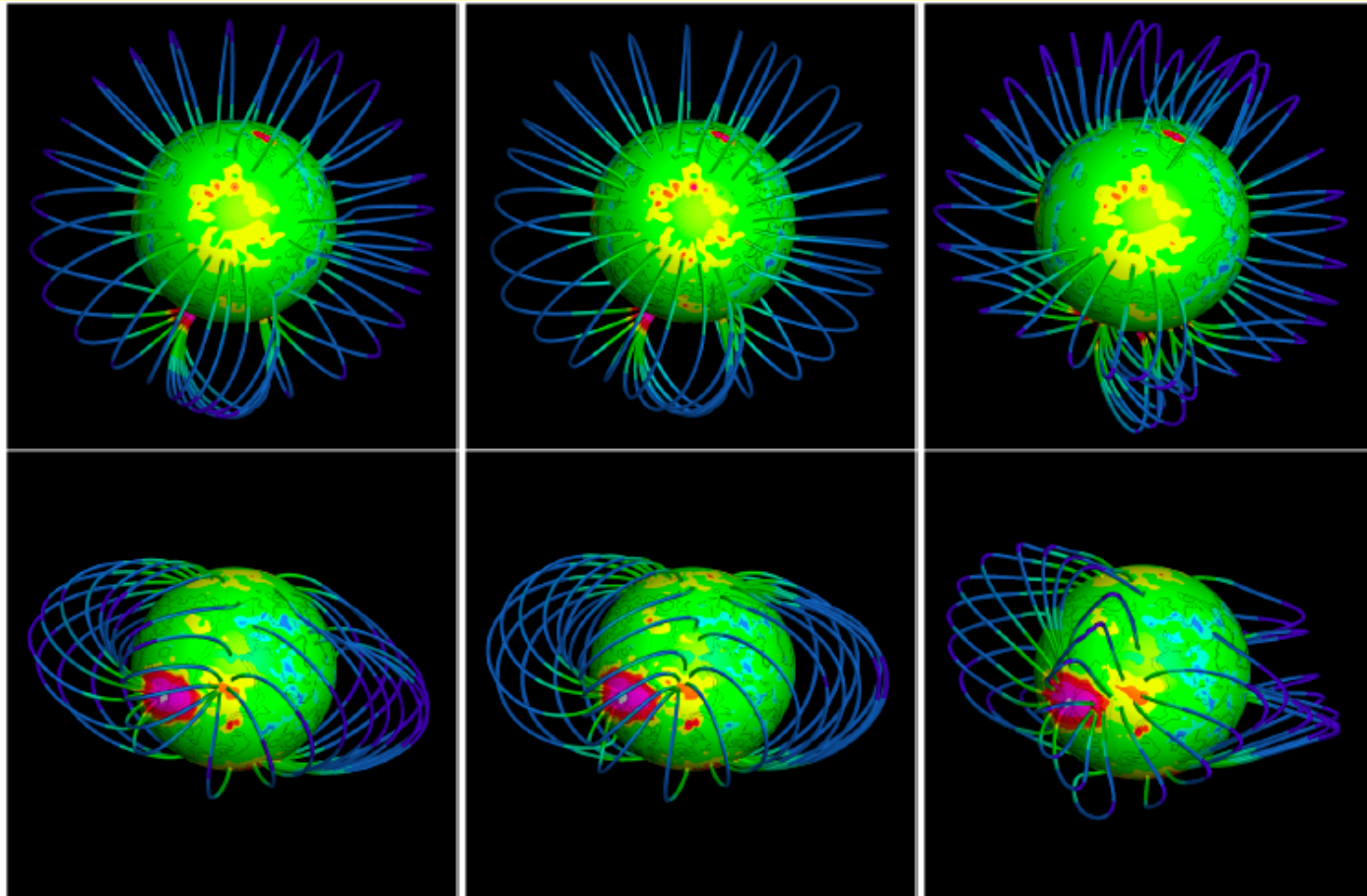
# Streamer structure for CR1913



PFSS model

Old SC model

New LC model



Top view

Earth  
Centered  
view

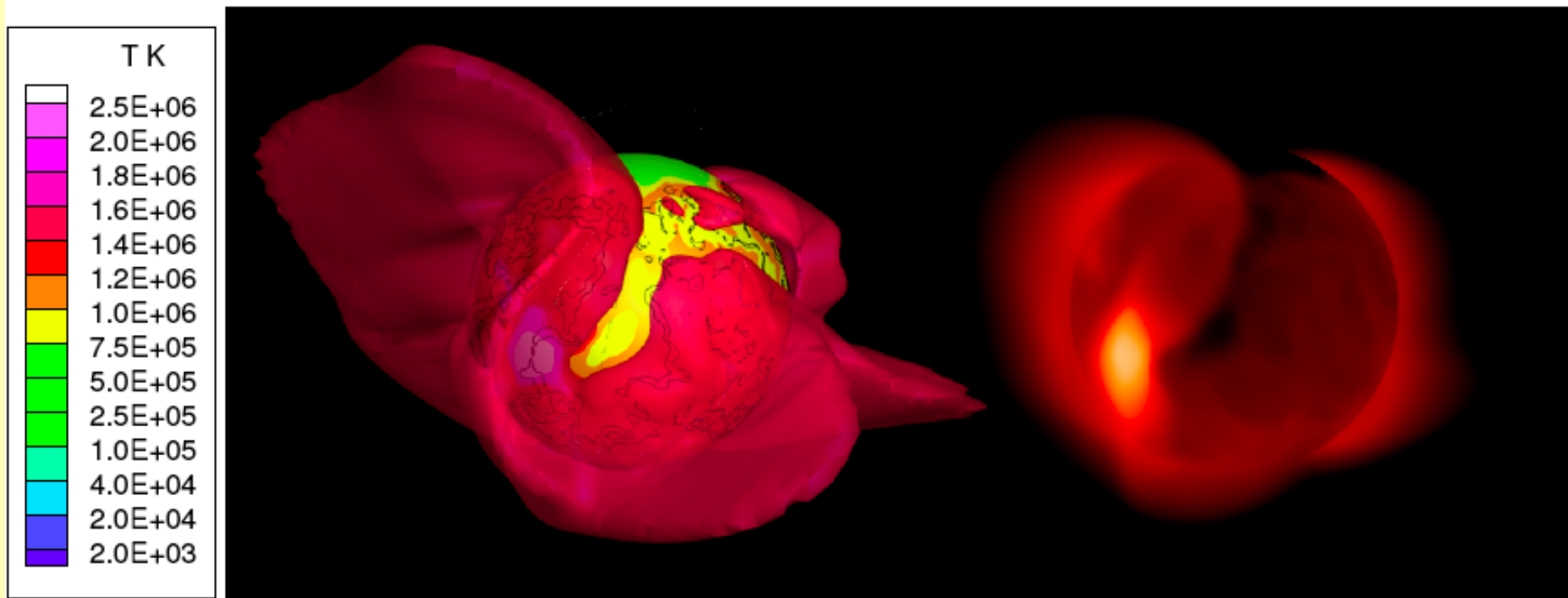
**M** Closed field lines structures are stressed by heating near surface and subsequent redistribution via heat conduction

# SXT response for CR1913



3D topology of T=1.6MK

LOS synthesis of SXT AlMg response



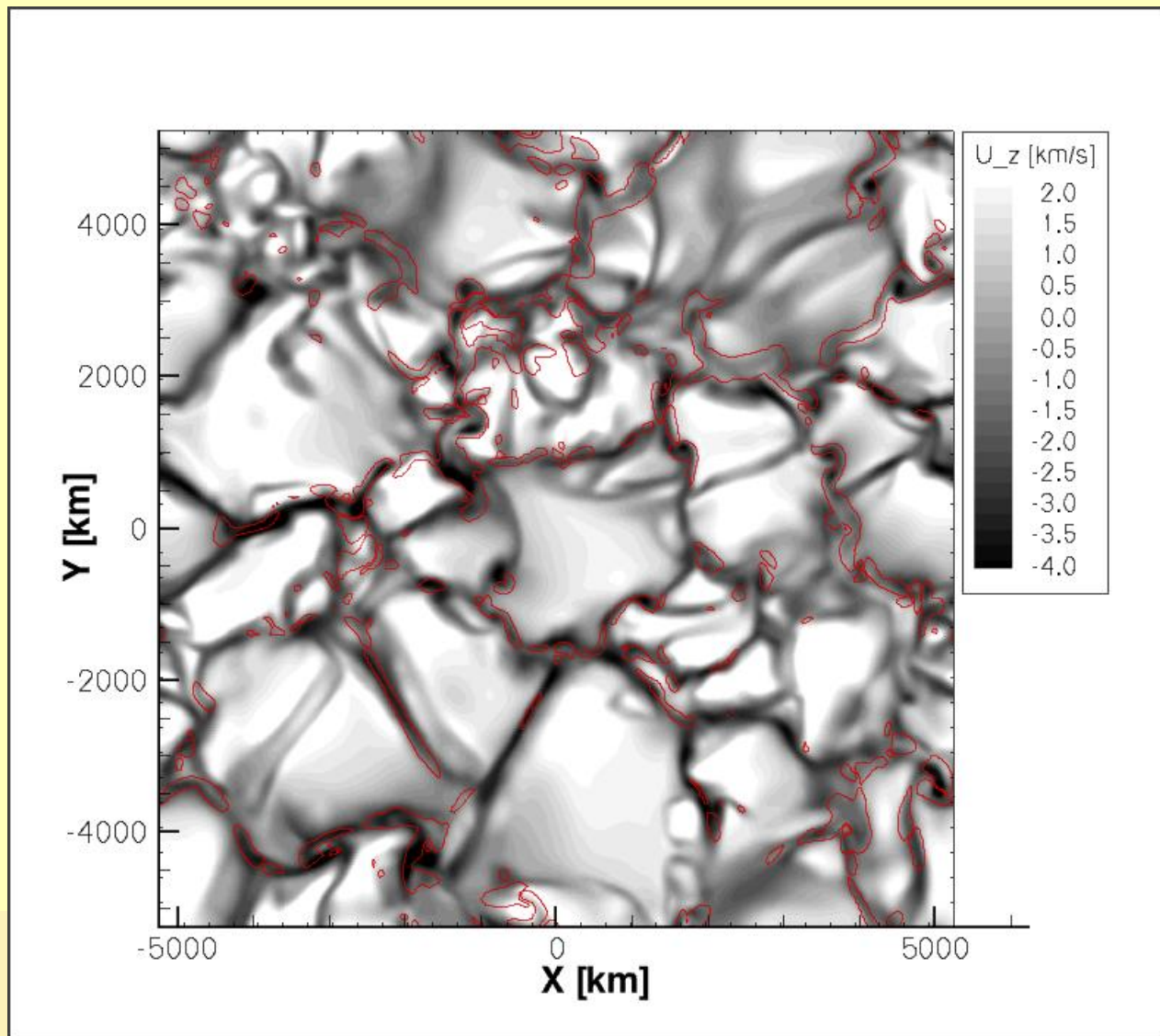
**M** The hot, closed field line regions overlay the inversion lines on the surface and correspond to the x-ray emission.



Atmosphere with coronal heating and radiative losses (Abbett, 2007)  
OPAL EOS

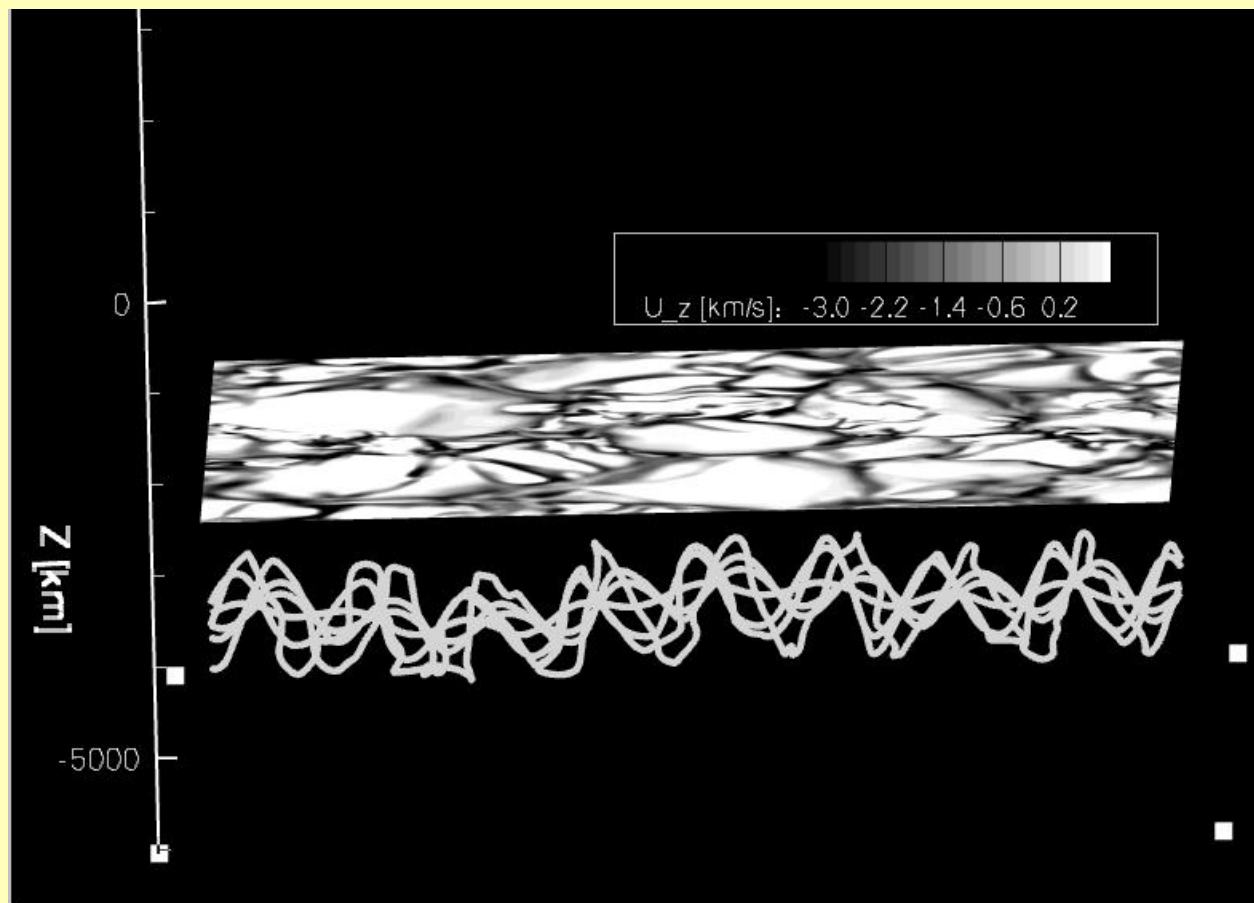
Vertical velocity at photospheric level shown in gray scale

Magnetic field concentrated in the intergranular lanes





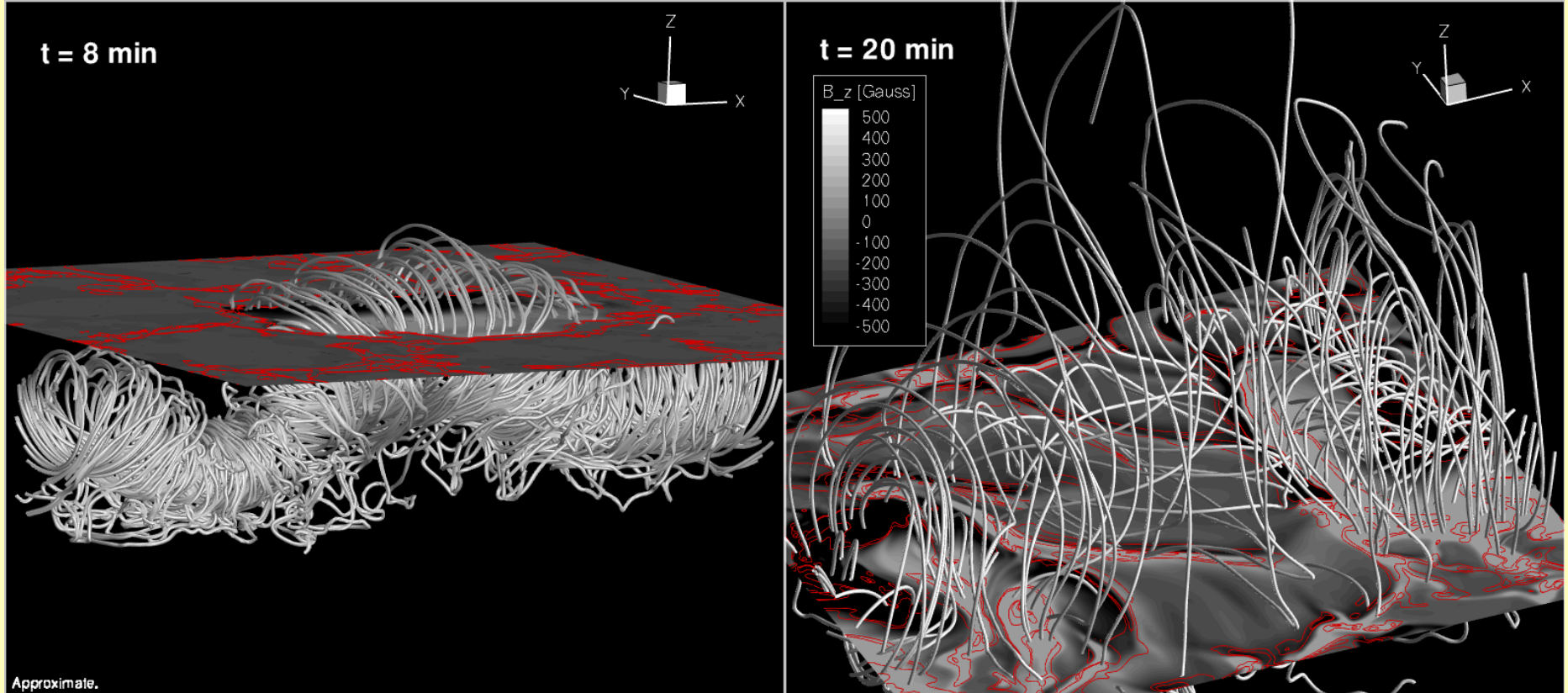
# Flux Emergence In the Corona



**M** Buoyant flux rope initially at 1500km below the photosphere

**M** 3D view at 2 minutes after initialization

# Flux Emergence In the Corona



**M** Convective down flow tries to drag flux rope downwards, up flow helps the flux to emerge, thus segmentation of rope

## Future plans



- M** Combine the two-temperature coronal model with the Alfvén wave turbulence model (see presentation Igor Sokolov)
- M** Use the newly developed radiation MHD package (gray and multigroup radiation diffusion) for the Convection Zone
- M** Couple the new Convection Zone, Lower Corona, Solar Corona, and Inner Heliosphere
- M** Do physically consistent CME eruptions into a physically consistent solar wind